

# One dimensional model systems in time-dependent density functional theory

N. Helbig, J.I. Fuks, I.V. Tokatly, and A. Rubio

Universidad del País Vasco, San Sebastián

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# Outline

- 1 1D model systems
- 2 Local density approximation
- 3 Rabi oscillations
- 4 Conclusions

Octopus code: [www.tddft.org/programs/octopus](http://www.tddft.org/programs/octopus)



# 1D model systems

## $N$ particles in 1D

$$\sum_{j=1}^N \left[ -\frac{d^2}{2dx_j^2} + v_{ext}(x_j) \right] + \frac{1}{2} \sum_{\substack{j,k=1 \\ j \neq k}}^N \frac{1}{\sqrt{(x_j - x_k)^2 + 1}}$$

- Equivalent to 1 particle in  $N$  dimensions
- Solve problem by **exact** diagonalization of Hamiltonian
- **Fermionic symmetry** enforced by projecting on Young diagrams

## Local density approximation

- Parametrization of the **correlation energy** from Quantum Monte-Carlo calculations of 1D homogeneous electron gas

# Linear response

Beryllium ions

$$v_{ext}(x, t) = -\frac{4}{\sqrt{x^2 + 1}} + \epsilon_0 \delta(t)$$

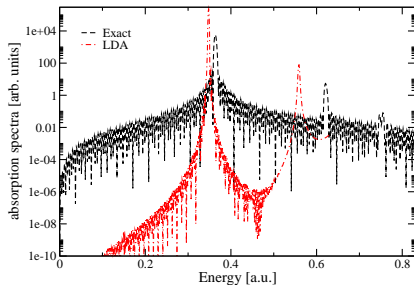
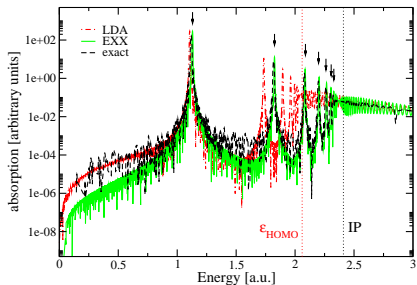


Fig. 1: Linear response spectra for Be<sup>2+</sup> (left) and Be<sup>+</sup> (right).

# Nonlinear response

Increased amplitude  $\mathcal{E}_0$

Second order response vanishes due to symmetry

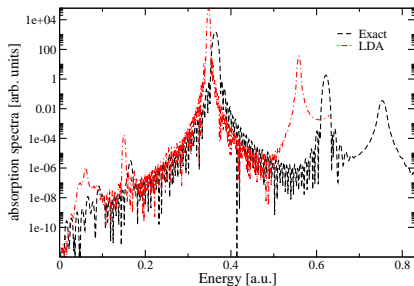
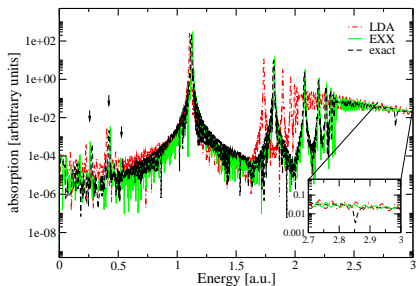


Fig. 2: Nonlinear response spectra for  $\text{Be}^{2+}$  (left) and  $\text{Be}^+$  (right).

## Rabi oscillations

$$v_{ext}(x, t) = -\frac{2}{\sqrt{x^2 + 1}} + \mathcal{E}_0 \sin((\omega_0 + \delta)t), \quad \omega_0 = 0.533 \text{ Ha}$$

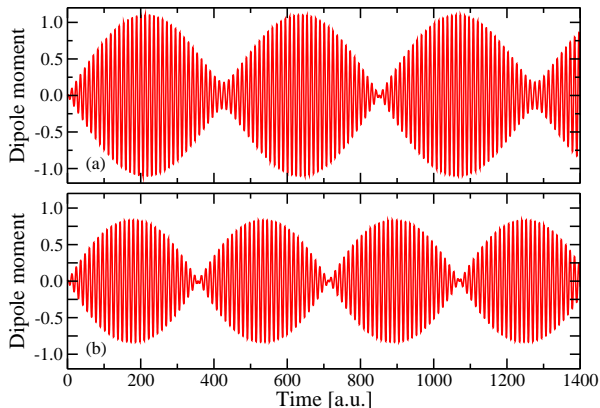


Fig. 3: Rabi oscillations with  $\delta = 0.0006$  Ha and  $\delta = 0.016$  Ha

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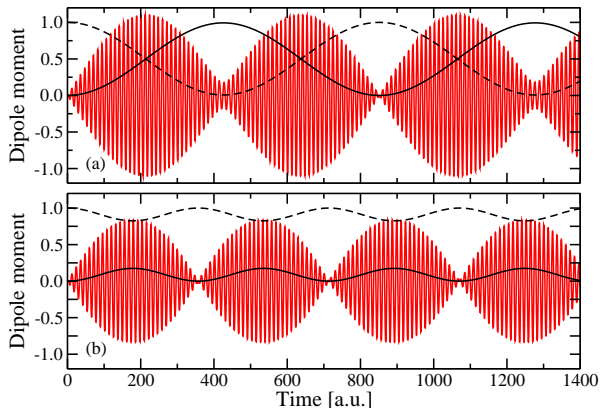


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# Rabi oscillations in TDDFT

Excitation with resonant frequency from linear response

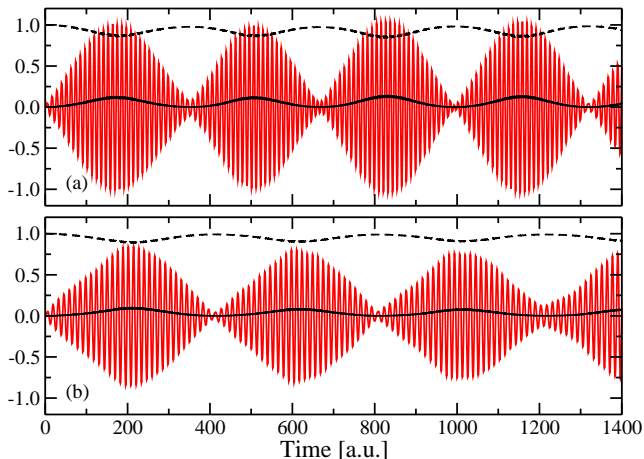


Fig. 4: Rabi oscillations in EXX (a) and LDA (b)



# Theoretical analysis

## Two level approximation

$$|\Psi(t)\rangle = a_0(t)|\Psi_0\rangle + a_1(t)|\Psi_1\rangle$$

$$n_j(t) = |a_j(t)|^2$$

- Exact calculation

$$i\partial_t \begin{pmatrix} a_0(t) \\ a_1(t) \end{pmatrix} = \begin{pmatrix} \epsilon_0 & d_{10}\mathcal{E}(t) \\ d_{10}\mathcal{E}(t) & \epsilon_1 \end{pmatrix} \begin{pmatrix} a_0(t) \\ a_1(t) \end{pmatrix}$$

- DFT calculation

$$i\partial_t \begin{pmatrix} a_0^s(t) \\ a_1^s(t) \end{pmatrix} = \begin{pmatrix} \epsilon_0^s + \epsilon_0^{xc}(t) & d_{10}^s\mathcal{E}(t) + \mathcal{F}_{xc}(t) \\ d_{10}^s\mathcal{E}(t) + \mathcal{F}_{xc}^*(t) & \epsilon_1^s + \epsilon_1^{xc}(t) \end{pmatrix} \begin{pmatrix} a_0^s(t) \\ a_1^s(t) \end{pmatrix}$$

# Theoretical model versus propagation

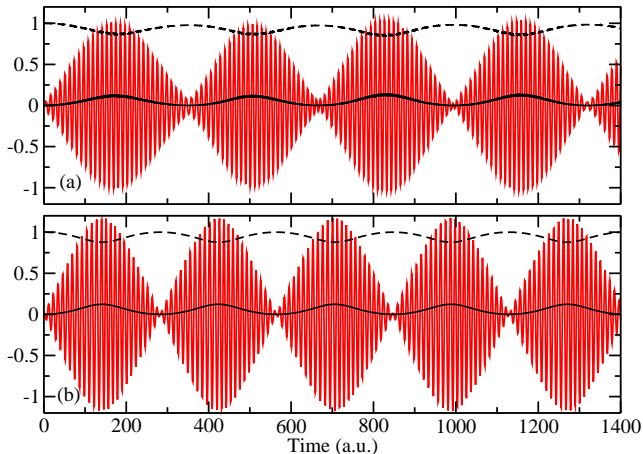


Fig. 5: Rabi oscillations from time propagation (a) and theoretical model (b) for the EXX functional

# Conclusions and Outlook

- 1D LDA of the same quality as its 3D counterpart
- Good approximation for **non-linear** response
- Adiabatic approximations lead to **detuning** in Rabi oscillations
- **Non-adiabatic** approximations

## References:

N. Helbig, J.I. Fuks, M. Casula, M.J. Verstraete, M.A.L. Marques, I.V. Tokatly, A. Rubio

[arXiv:1101.2564](https://arxiv.org/abs/1101.2564) accepted in PRA

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