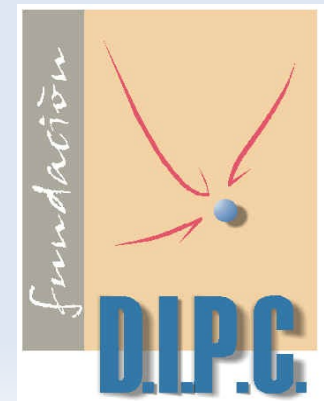


Non-adiabatic effects in one-dimensional one- and two-electron systems: the cases of H_2^+ and H_2

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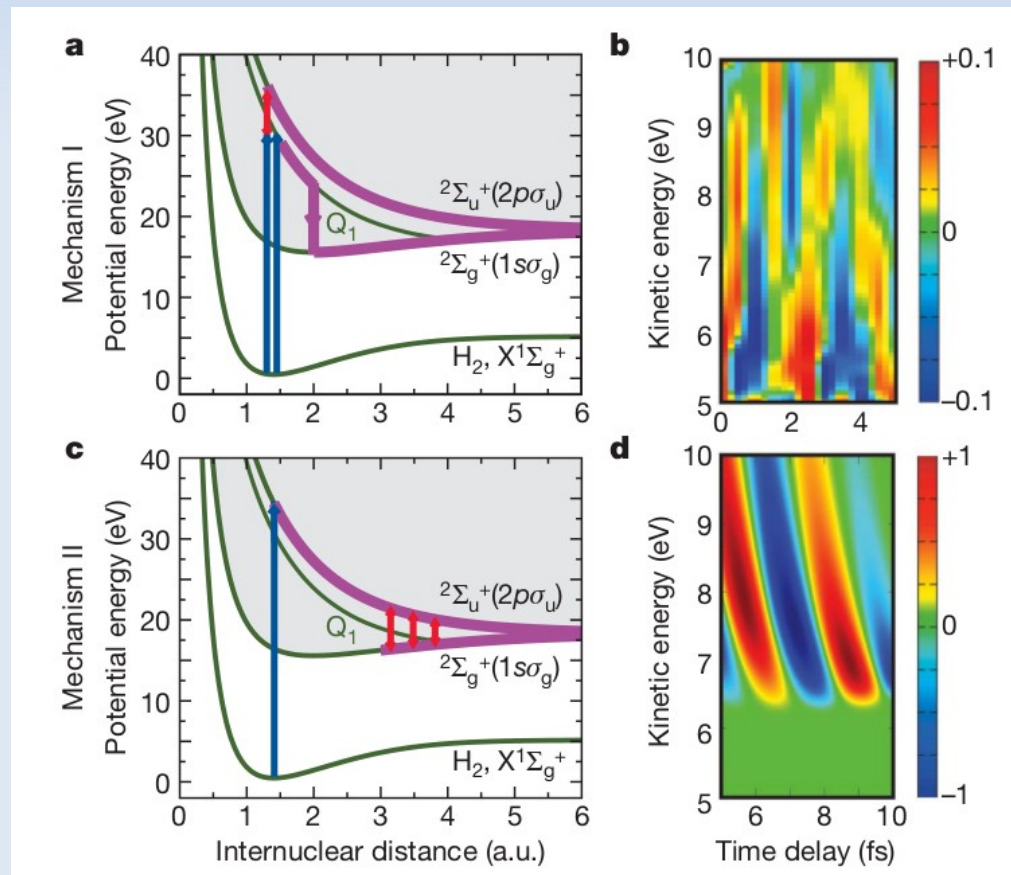


Outline

- Motivations
- Model systems: one-dimensional H_2^+ and H_2
- Results:
 - Validity of the Born Oppenheimer Approximation
 - Optical spectra from frozen ion calculations
 - Optical spectra from dynamic ion calculations
- Conclusions
- Future work

Motivations

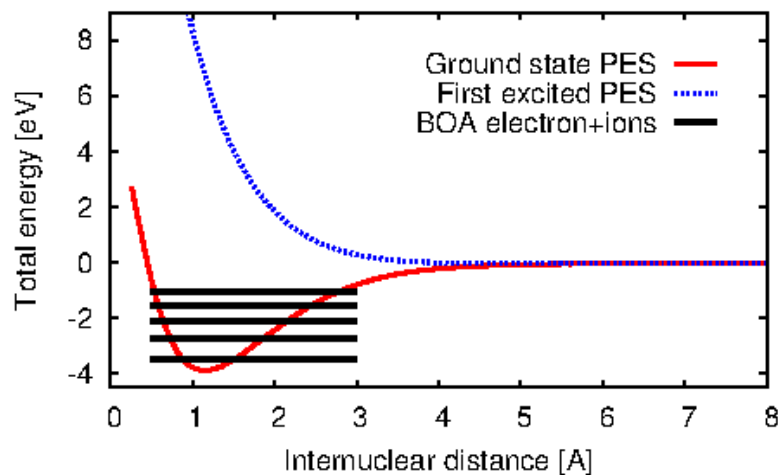
- Interpret pump-probe attosecond experiments beyond the Born Oppenheimer Approximation



The Born Oppenheimer Approximation

- Assess the validity of the Born-Oppenheimer Approximation (BOA)

$$\Psi_{total} \approx \Psi_{BOA} = \Psi_{electronic} \Psi_{ionic}$$



Potential Energy Surfaces (PES's)

- If $\frac{m_e}{m_I} \ll 1$, the kinetic energy of the ions is negligible: “frozen ions”
- Fictitiously vary the electron-ion mass ratio $\frac{m_e}{m_I}$ to change the “electron-ion coupling”

S. Takahashi and K. Takatsuka, J. Chem. Phys. 124 (2006), 1–14.

Model systems: H_2^+ and H_2 in 1D

- The exact numerical diagonalisation in real-space is feasible
- Exchange symmetry of the molecular wavefunction
The spin part is directly determined (singlet, triplet)

$$\psi_{H_2^+}(R_1 S_1, R_2 S_2, r_1 s_1) = -\psi_{H_2^+}(R_2 S_2, R_1 S_1, r_1 s_1)$$

$$\psi_{H_2}(R_1 S_1, R_2 S_2, r_1 s_1, r_2 s_2) = -\psi_{H_2}(R_2 S_2, R_1 S_1, r_1 s_1, r_2 s_2)$$

$$\psi_{H_2}(R_1 S_1, R_2 S_2, r_1 s_1, r_2 s_2) = -\psi_{H_2}(R_1 S_1, R_2 S_2, r_2 s_2, r_1 s_1)$$

- Soft Coulomb Potential

Coulomb potential ill-defined in 1-D

R. Loudon, Am J. Phys. 27 (1959), 649-655

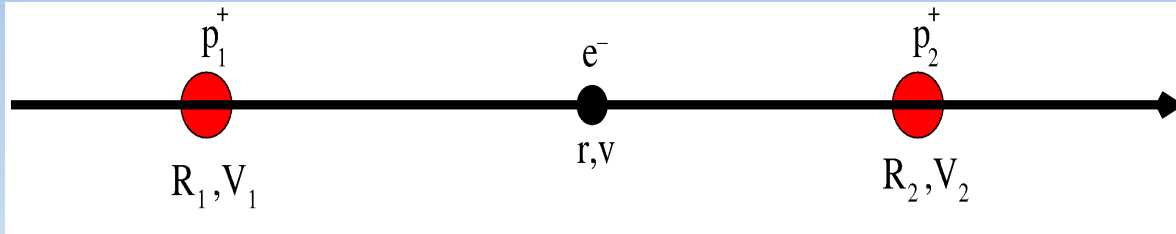
$$V_{int}(x_i - x_j) = \frac{q_i q_j}{\sqrt{(x_i - x_j)^2 + a^2}}$$

- We use the real-space code OCTOPUS

A. Castro et al., phys. stat. Sol. 243 (2006), 2465-2488

http://www.tddft.org/programs/octopus/wiki/index.php/Main_page

The 1D dihydrogen cation H_2^+



$$R = R_2 - R_1$$

$$\xi = r - \frac{R_1 + R_2}{2}$$

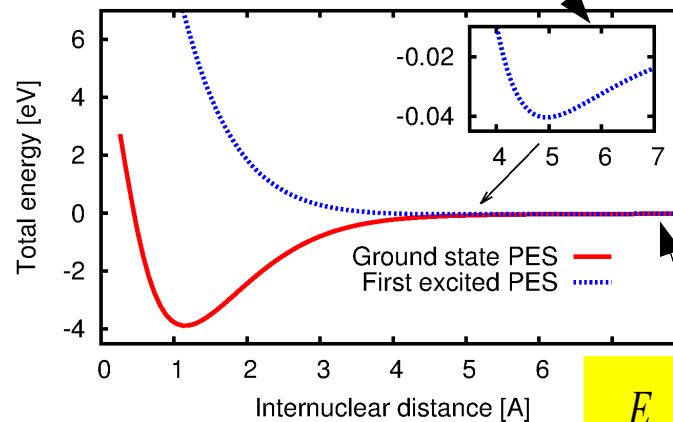
- Hamiltonian (centre of mass frame) in atomic units (a.u.)

J. R. Hiskes, Phys. Rev. 122 (1960), 1207-1217

$$H_{internal}(R, \xi) = \frac{1}{2\mu_I} \frac{\partial^2}{\partial R^2} - \frac{1}{2\mu_e} \frac{\partial^2}{\partial \xi^2} - \frac{1}{\sqrt{\left(\frac{R}{2} + \xi\right)^2 + 1}} - \frac{1}{\sqrt{\left(\frac{R}{2} - \xi\right)^2 + 1}} + \frac{1}{\sqrt{R^2 + 1}}$$

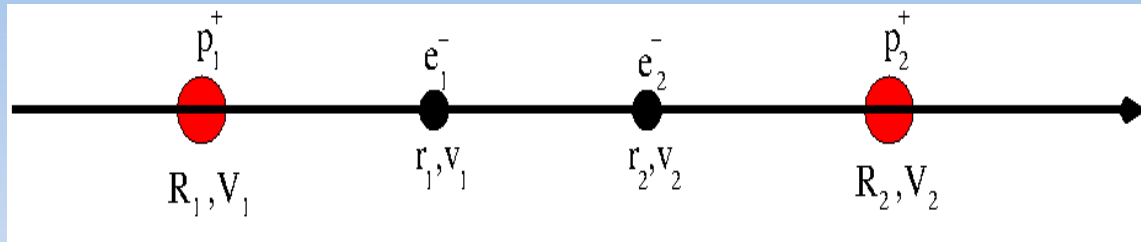
Negligible if $\mu_I \gg \mu_e$

Non-covalent long range minimum ($H^+ - H$)



$$E_{gs}(R) \sim -\frac{1}{R^3}$$

The 1D dihydrogen H₂



$$R = R_2 - R_1 \quad r = r_2 - r_1$$

$$\xi = \frac{r_1 + r_2}{2} - \frac{R_2 + R_1}{2}$$

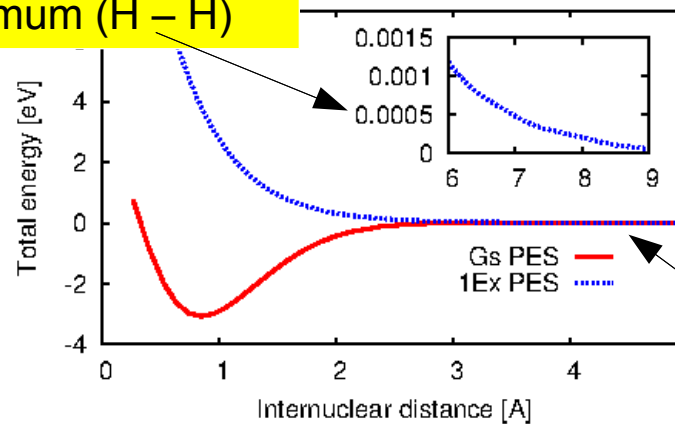
- Hamiltonian (centre of mass frame) in atomic units (a.u.)

$$H_{internal}(R, r, \xi) = \frac{1}{2\mu_I} \frac{\partial^2}{\partial R^2} - \frac{1}{2\mu_{eI}} \frac{\partial^2}{\partial \xi^2} - \frac{1}{2\tilde{\mu}_e} \frac{\partial^2}{\partial r^2} - \frac{1}{\sqrt{\left(\frac{R}{2} - \frac{r}{2} + \xi\right)^2 + 1}} - \frac{1}{\sqrt{\left(\frac{R}{2} - \frac{r}{2} - \xi\right)^2 + 1}}$$

Negligible if $\mu_I \gg \tilde{\mu}_e$

$$- \frac{1}{\sqrt{\left(\frac{R}{2} + \frac{r}{2} + \xi\right)^2 + 1}} - \frac{1}{\sqrt{\left(\frac{R}{2} + \frac{r}{2} - \xi\right)^2 + 1}} + \frac{1}{\sqrt{R^2 + 1}} + \frac{1}{\sqrt{r^2 + 1}}$$

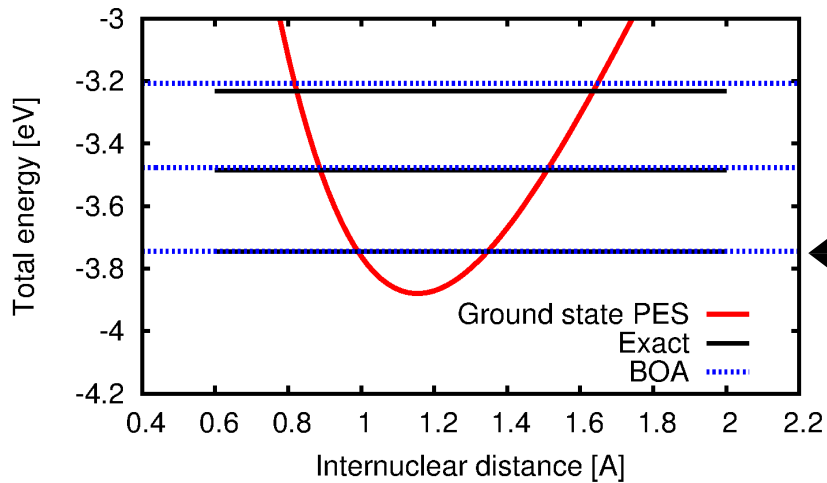
Non-covalent long range minimum (H - H)



$$E_{gs}(R) \sim -\frac{1}{R^3}$$

BOA validity: H_2^+ case

Bottom of the ground state PES



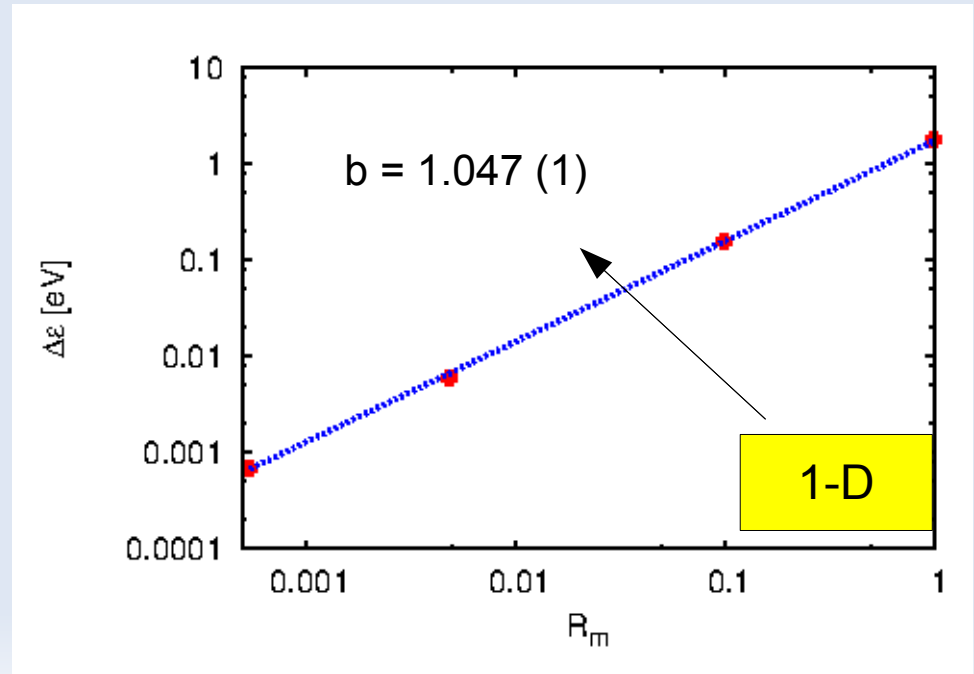
$\frac{m_e}{m_I}$	m_I	$E_{EXACT} [eV]$	$E_{BOA} [eV]$	$\Delta E [eV]$
5.45×10^{-4}	(proton)	-3.7454(3)	-3.7447(7)	0.0007(5)
4.84×10^{-3}	(muon)	-3.4851(3)	-3.4791(4)	0.0060(4)
0.1	(10 electron)	-2.2525(2)	-2.0936(9)	0.1589(6)
1.0	(electron)	-0.6052(2)	1.165(3)	1.770(3)

E_{BOA} = bottom PES + zero-point energy $\left(\frac{1}{2} \hbar \omega\right)$
 E_{EXACT} (numerical)

- BOA, expansion E_{gs} in terms $\left(\frac{m_e}{m_I}\right)^{\frac{1}{4}}$

$$\Delta E = E_{BOA} - E_{EXACT} = a \left(\frac{m_e}{m_I}\right)^b$$

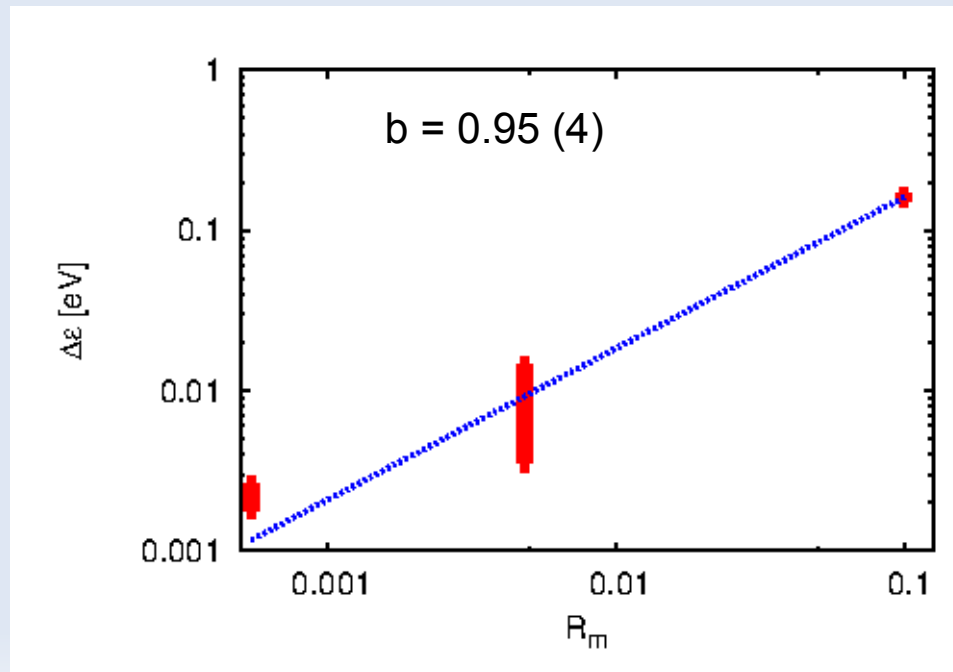
- 3-D $\rightarrow b=1.5$
- 1-D $\rightarrow b=1$ (There are no contributions from rotations)



BOA validity: H₂ case

$\frac{m_e}{m_I}$	m_I	$E_{EXACT} [eV]$	$E_{BOA} [eV]$	$\Delta E [eV]$
5.45×10^{-4}	(proton)	-2.89071(3)	-2.8885(5)	0.0022(3)
4.84×10^{-3}	(muon)	-2.53937(2)	-2.530(1)	0.009(1)
0.1	(10 electron)	-1.11145(2)	-0.94962(3)	0.16183(3)

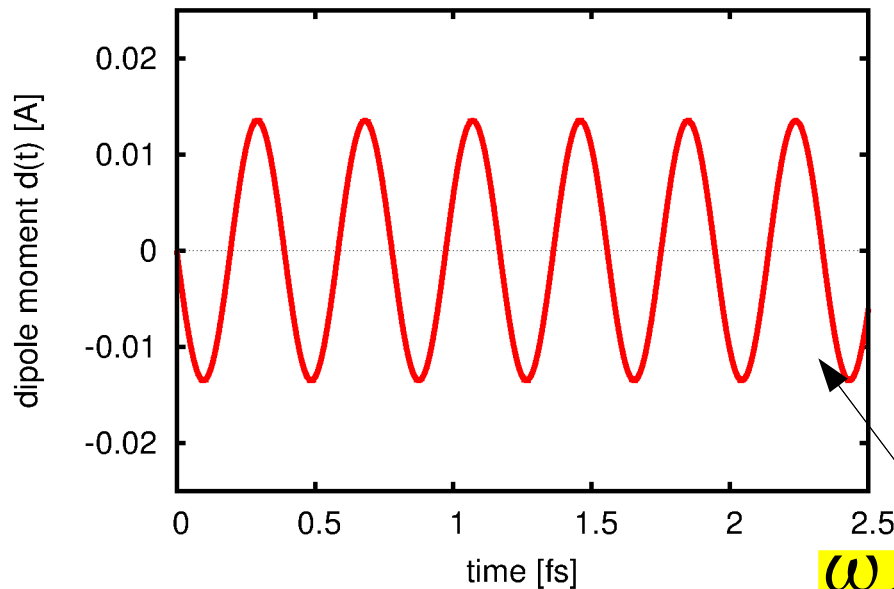
- No bound states for $\frac{m_e}{m_I} = 1$ (**electron**)



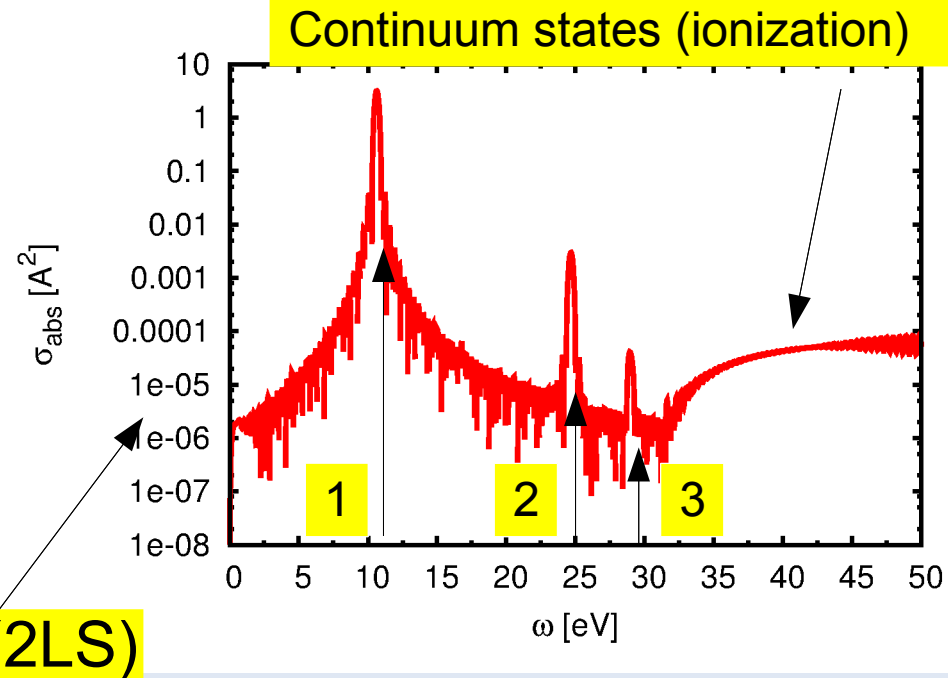
Frozen Ion Optical Spectra H_2^+

- The system is perturbed by a weak "kick" $|\psi(r, t=0)\rangle = e^{ikr} |\psi_{gs}\rangle$
- Dipole response $d(t)$

$$d(t) = -\bar{d} \sin(\omega_{eq} t)$$



$$\sigma_{abs} = 4\pi\omega\alpha \text{Im}\left[\frac{1}{k} \int_0^t dt e^{-i\omega t} f\left(\frac{t}{T}\right) d(t)\right]$$



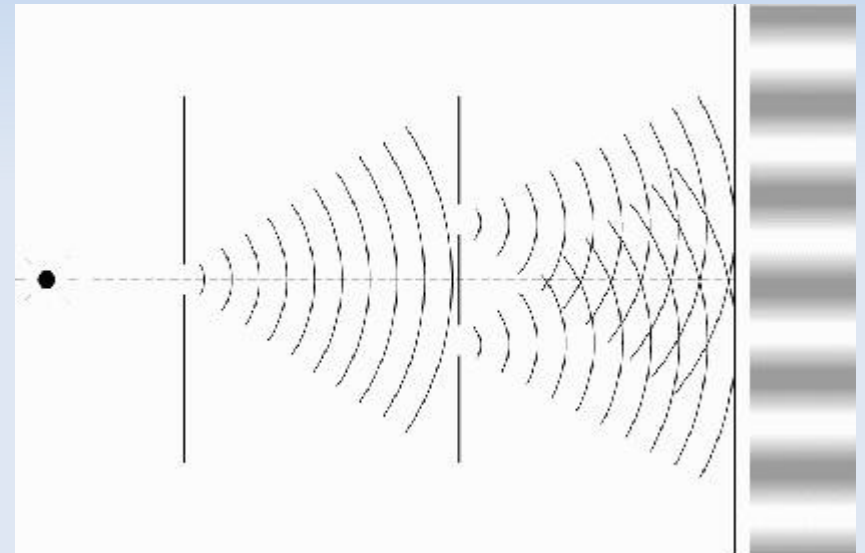
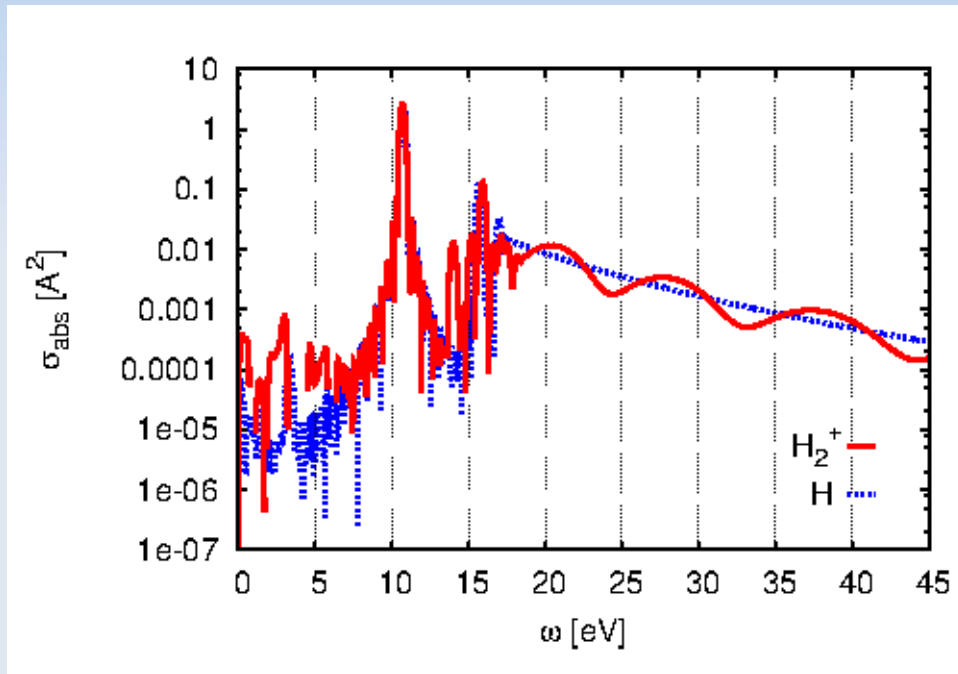
1: Ground State \rightarrow First Excited State (ω_{eq})

2: Ground State \rightarrow Third Excited State

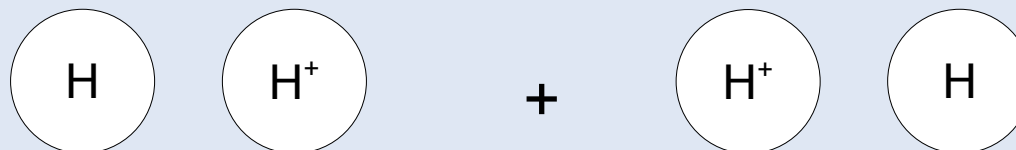
3: Ground State \rightarrow Fifth Excited State

Frozen Ion Optical Spectra H_2^+

- For large R: H_2^+ vs H



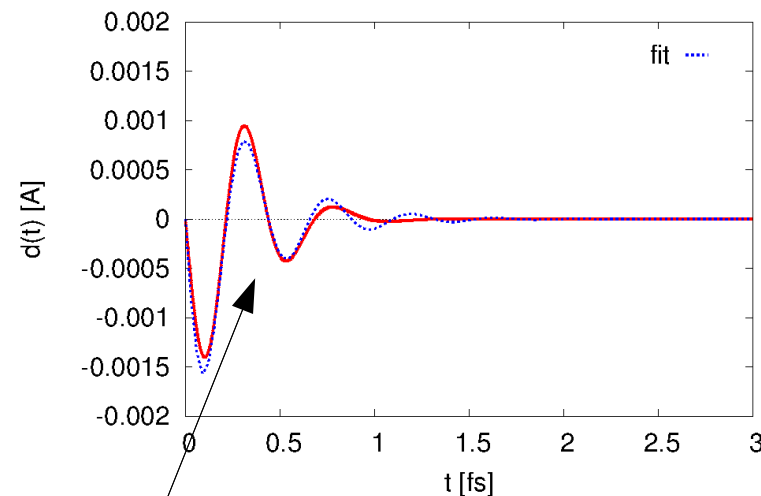
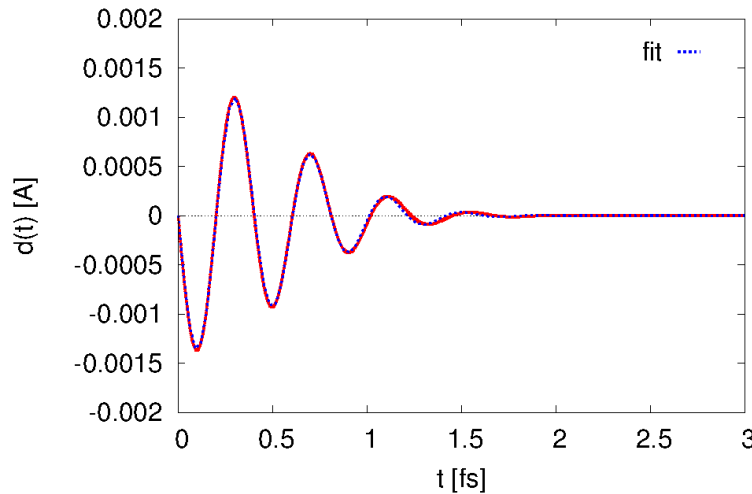
<http://www.physics.uiowa.edu/~umallik/adventure/quantumwave.html>



Dynamic Ion Optical Spectra H_2^+

$$\frac{m_e}{m_I}$$

$$d(t) = -\bar{d} e^{-\frac{b^2 t^2}{2}} \sin(\bar{\omega} t)$$

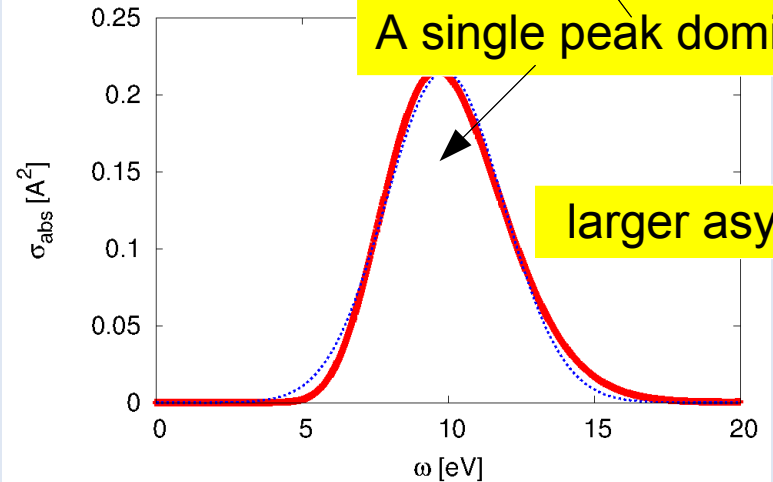
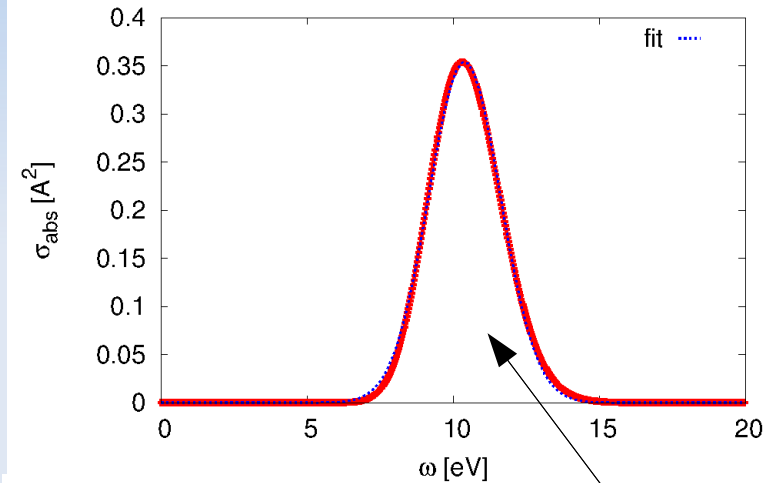


Quicker energy transfer

5.45×10^{-4}
(proton)

4.84×10^{-3}
(muon)

$$\sigma_{abs} = \frac{a \omega}{\sqrt{2 \pi b^2}} e^{-\frac{(\omega - \omega_0)^2}{2b^2}}$$

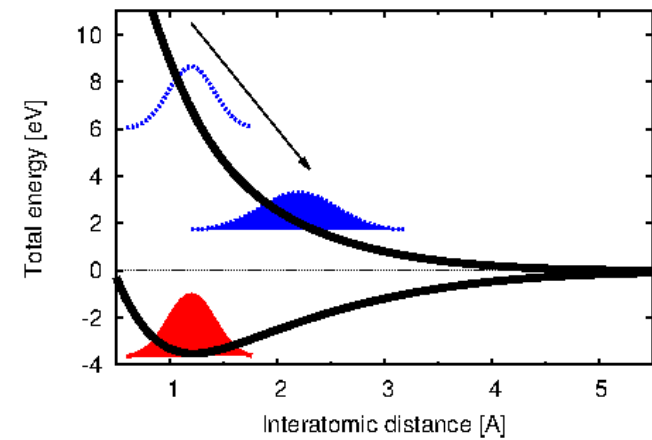
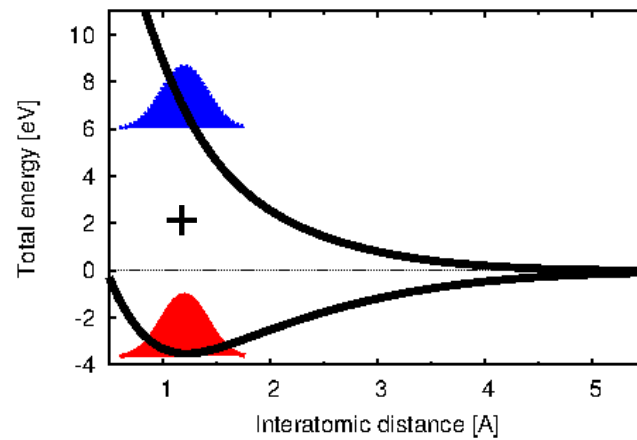
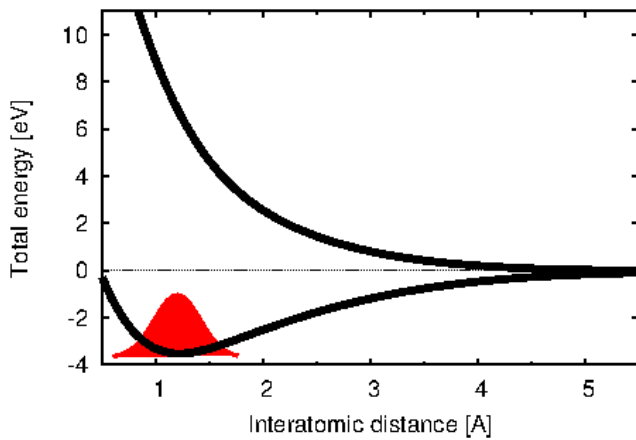


A single peak dominates

larger asymmetry

Dynamic Ion Optical Spectra H_2^+

- Gaussian qualitative analysis (2LS)



$$d(t) = -\bar{d} e^{-\frac{b^2 t^2}{2}} \sin(\bar{\omega} t)$$

$$\sigma_{abs} = \frac{a \omega}{\sqrt{2 \pi b^2}} e^{-\frac{(\omega - \omega_0)^2}{2b^2}}$$

$\frac{m_e}{m_I}$	m_I	$\hbar b [eV]$	$\hbar \bar{\omega} [eV]$	$b [eV]$	$\omega_0 [eV]$
5.45×10^{-4}	(proton)	1.052(2)	10.228(7)	1.237(3)	10.237(3)
4.84×10^{-3}	(muon)	2.02(1)	9.428(6)	2.051(9)	9.44(1)

Conclusions

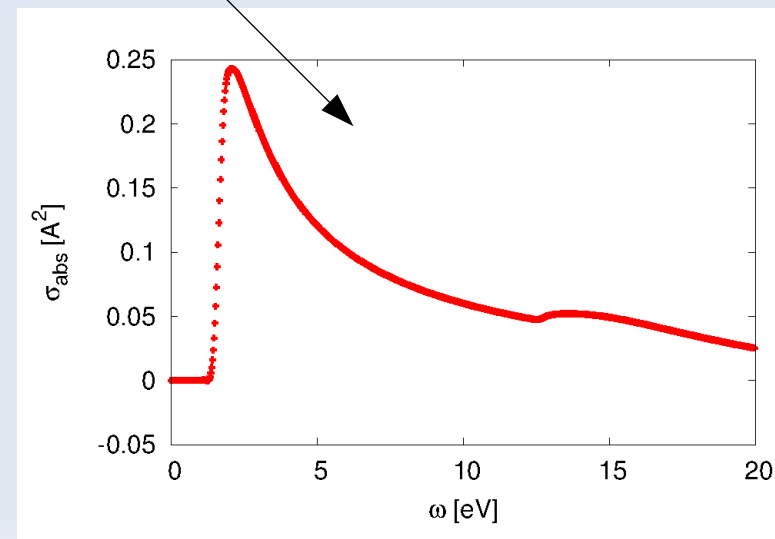
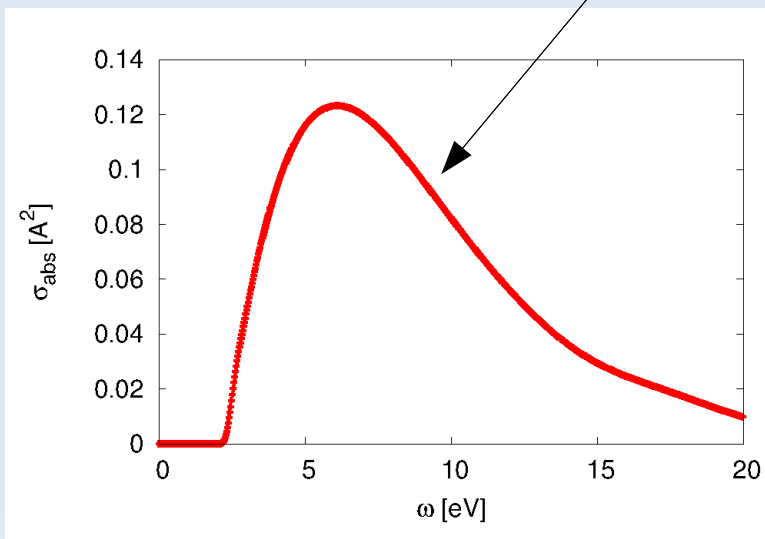
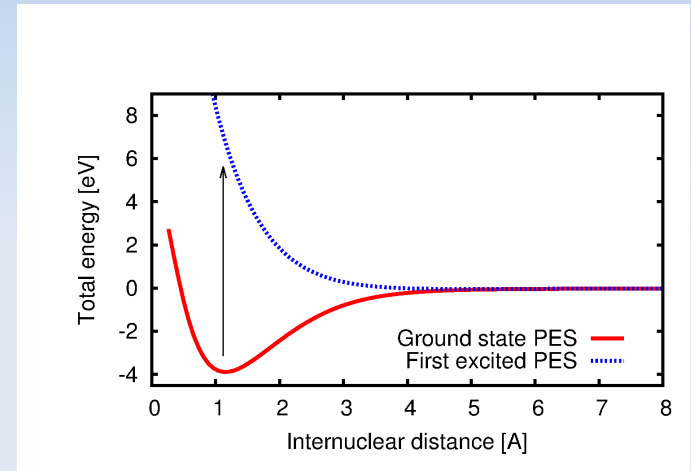
- Static case H_2^+ and H_2 , we find $b \approx 1$ (1-D)

$$E_{BOA} - E_{EXACT} = a \left(\frac{m_e}{m_I} \right)^b$$

- Dynamic case H_2^+ , single frequency \rightarrow two-level system (2LS)

dynamics for small $\frac{m_e}{m_I}$ (proton, muon)

- 2LS is not accurate for $\frac{m_e}{m_I} = 0.1$; $\frac{m_e}{m_I} = 1$ (yet to be fully understood)



Future Work

- Dihydrogen H_2 optical spectra from frozen and dynamic ion calculations
- Improve theoretical model for the dynamic ion calculations (asymmetry, gaussian)
- Perform TDDFT and Ehrenfest dynamics calculations and compare to the exact calculations
- Consider more realistic systems and electromagnetic pulses (pulse shapes) to interpret the experiments

Non-adiabatic effects in one-dimensional one and two electron systems:

the cases of H_2^+ and H_2

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THANK YOU