

# Non-adiabatic effects in the optical spectra of $H_2^+$

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# Outline

- Motivations:
  - Ions are usually treated classically (BOA)
  - Effects of a full quantum treatment ?
- Model system:  $H_2^+$  in 1D
- Optical spectra: BOA vs full quantum
- BOA theoretical analysis
- Conclusions



# Motivations

- Full quantum treatment

$$H_{eI}(R, r) \psi_{eI}(R, r) = E_{eI}(R, r) \psi_{eI}(R, r)$$

$$H_{eI}(R, r) = T_I(R) + T_e(r) + V_{eI}(R, r) + V_{ee}(r) + V_{II}(R)$$

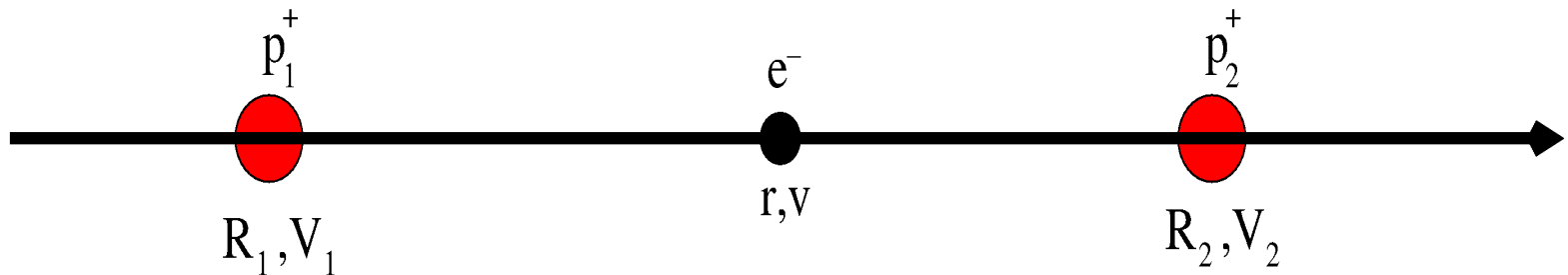
- Adiabatic approximations  $\frac{m_e}{m_I} \ll 1 \rightarrow$  decouple electron-ion
  - Born Oppenheimer Approximation (BOA)

$$\psi_{eI}(R, r) \approx \psi_{BOA}(R, r) = \chi_I(R) \varphi_e^R(r)$$

- Electron:  $\varphi_e^R(r) H_e = \varepsilon(R) \varphi_e^R(r); H_e = T_e + V_{eI} + V_{ee} + V_{II}$
- Ion:  $H_I^i \chi_I(R, t) = E_I^i(R) \chi_I(R, t); H_I^i = T_I + \varepsilon_i(R)$

**See effects of BOA vs full quantum treatment**

# Model system: $H_2^+$ in 1D



- Hamiltonian in 1D (centre of mass)

$$H_{internal}(R, \xi) = -\frac{1}{2\mu_I} \frac{\partial^2}{\partial R^2} - \frac{1}{2\mu_e} \frac{\partial^2}{\partial \xi^2} - \frac{1}{\sqrt{\left(\frac{R}{2} + \xi\right)^2 + 1}} - \frac{1}{\sqrt{\left(\frac{R}{2} - \xi\right)^2 + 1}} + \frac{1}{\sqrt{R^2 + 1}}$$

Negligible if  $\mu_I \gg \mu_e$

$$R = R_2 - R_1 \quad \xi = r - \frac{R_1 + R_2}{2} \quad \mu_I = \frac{m_I}{2} \quad \mu_e = \frac{2m_I}{2m_I + m_e}$$

- Exact numerical diagonalisation feasible
- Octopus code

Change only  $m_I \rightarrow H_2^+ \dots D_2^+ \dots Li_2^+ \dots K_2^+ \rightarrow$  assess electron-ion coupling



# Optical spectra → theory

- Kicked initial state in the dipole approximation

$$\psi_{eI}(R, r) = \psi_0(R, r) e^{ikr} \approx (1 + ikr) \psi_0(R, r)$$

$$\psi_{eI}(R, r, t) \approx \psi_0(R, r) e^{\frac{H_0 t}{i\hbar}} + ik \sum_{i>0} \psi_i^*(R, r) r \psi_0(R, r) \psi_i(R, r) e^{\frac{H_i t}{i\hbar}}$$

- Dipole moment

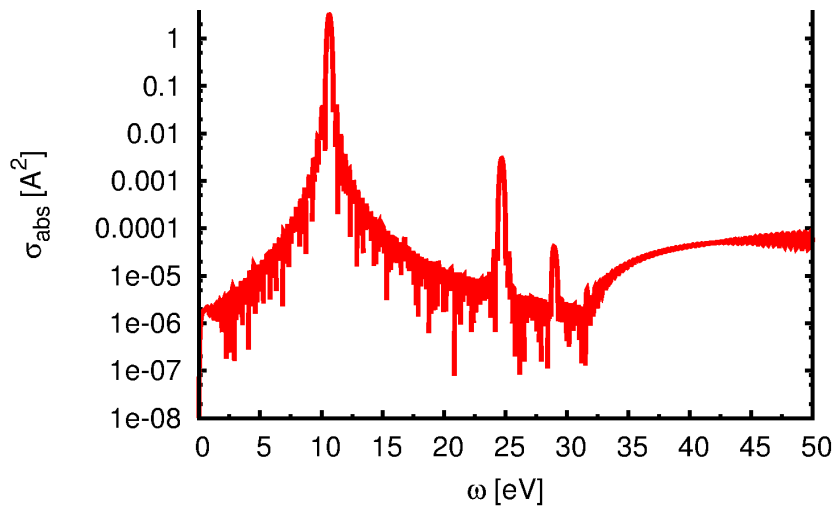
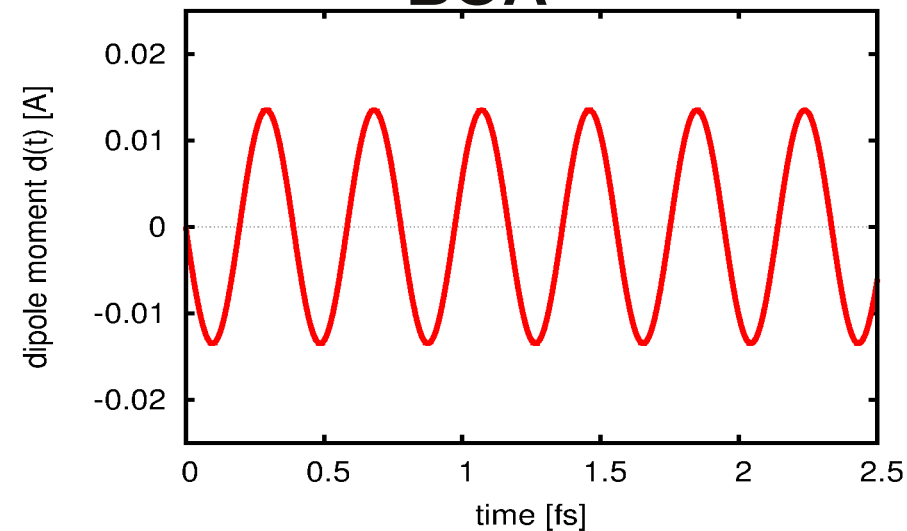
$$d(t) = \int \psi_{eI}^*(R, r, t) r \psi_{eI}(R, r, t)$$

- Optical spectra

$$\sigma(\omega) = \lim_{k \rightarrow \infty} \frac{4\pi\alpha\omega}{k} \Im \left[ \int_0^{\infty} dt e^{-i\omega t} (d(t) - d(0)) \right]$$

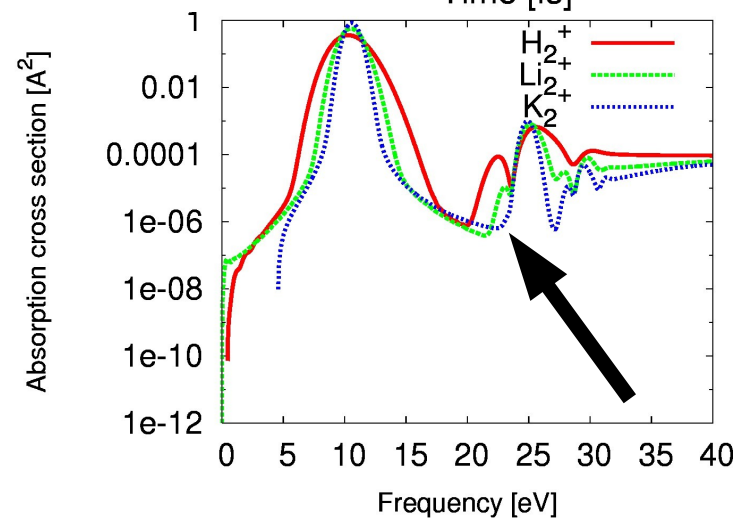
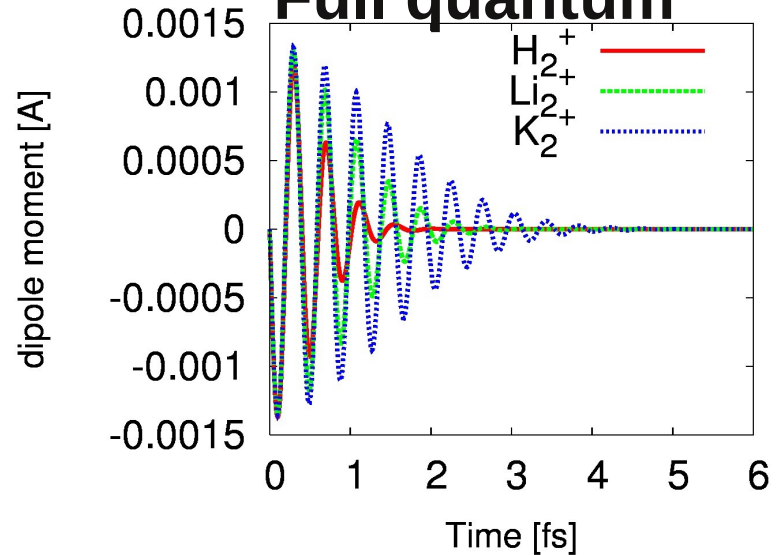
# Optical spectra: BOA vs Full

## BOA



**Lorentzian**  
 $m_I$  no peak shift

## Full quantum



**Gaussian**  
 $m_I$  peak shift  
**NEW PEAK !**

# BOA theoretical analysis

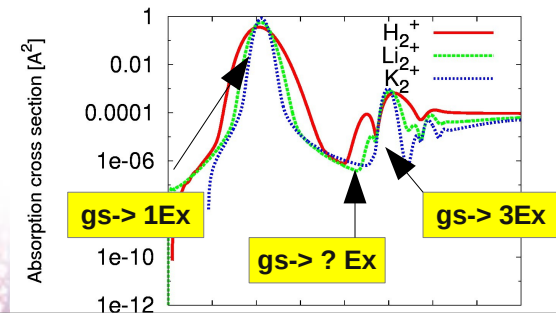
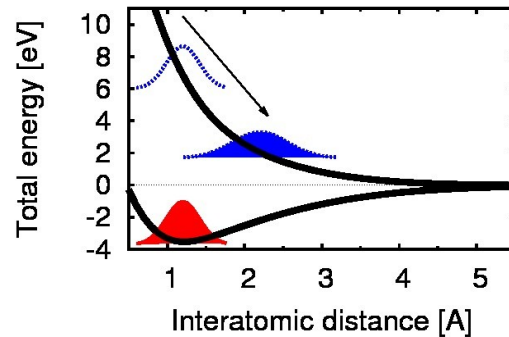
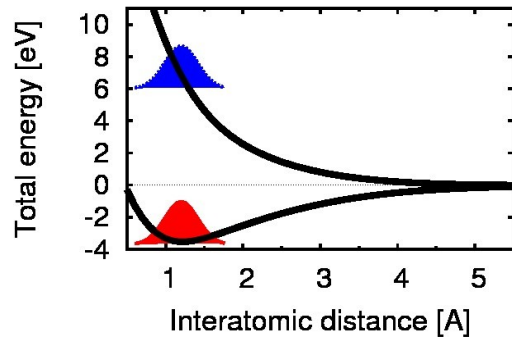
$$\Psi_{BOA}(R, r, t) \approx \chi_0(R) \varphi_0^{(R)}(r) e^{\frac{H_e t}{i\hbar}} + ik \sum_{i>0} d_{i0} \varphi_i^{(R)}(r) \chi_i(R, t)$$

$$\chi_i(R, t) = N e^{-\alpha_i(R-R_i)^2 - \frac{1}{i\hbar}(R-R_i)P_i - \frac{1}{i\hbar}\gamma_i}$$

$$\chi_0(R, t) = N e^{-\alpha_0(R-R_0)^2} e^{\frac{E_0 t}{i\hbar}}$$

$$i\hbar \frac{\partial}{\partial t} \chi_i(R, t) = \left( -\frac{\hbar^2}{2\mu_I} \frac{\partial^2}{\partial R^2} - F_i(R-R_i) + \varepsilon_i(R_i) \right) \chi_i(R, t)$$

$$E_0 \chi_0(R) = \left( -\frac{\hbar^2}{2\mu_I} \frac{\partial^2}{\partial R^2} - \frac{1}{2} K (R-R_0)^2 + \varepsilon_0(R_0) \right) \chi_0(R)$$



$$d(t) = \int \Psi_{eI}^*(R, r, t) r \Psi_{eI}(R, r, t)$$



# BOA theoretical analysis

$$d(t) = \int \psi_{BOA}^*(R, r, t) r \psi_{BOA}(R, r, t)$$

$$d(t) = 2k \sum_{i>0} d_{i0}^2 \mathfrak{I} \int dR e^{-\frac{E_0 t}{i \hbar}} \chi_0^*(R) \chi_i(R, t)$$

$$d(t) = 2k \sum_{i>0} d_{i0}^2 \mathfrak{I} \left( e^{-\frac{F_i^2 \sigma_0^2}{2 \hbar^2} t^2 - \omega_i t + \frac{F_i^2}{12 \hbar \mu_I} t^3} \right)$$

**ONLY ADIABATIC EFFECTS**



# BOA theoretical analysis

$$\psi_{eI}(R, r, t) = \sum_i c_i(t) \chi_i(R, t) \varphi_i^R(r)$$

$$\langle \chi_n(R, t) \varphi_n(R, r) | i \hbar \frac{\partial \psi}{\partial t} - H \psi \rangle = 0$$

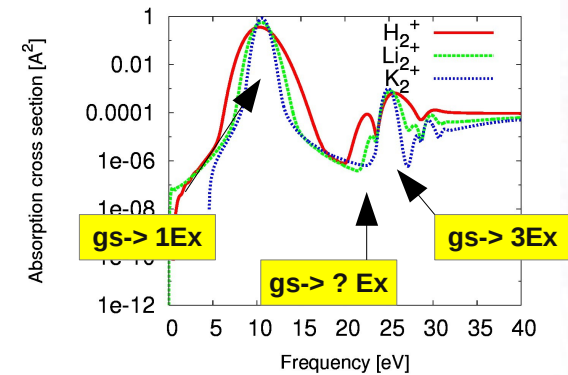
~~$$i \hbar \dot{c}_n = \frac{P_n}{\mu_I} \sum_i I_{ni} A_{ni} c_i + \frac{2i \hbar \alpha_0}{\mu_I} \sum_i I_{ni}^{(i)} A_{ni} c_i - \frac{1}{2\mu_I} \sum_i I_{ni} B_{ni} c_i$$~~

$$I_{ni} = \int dR \chi_n^*(R, t) \chi_i(R, t) \quad I_{ni}^{(i)} = \int dR \chi_n^*(R, t) (R - R_i) \chi_i(R, t)$$

$$A_{ni} = -i \hbar \langle \varphi_j(R) | \frac{\partial \varphi_i(R)}{\partial R} \rangle \quad B_{ni} = -\hbar^2 \langle \varphi_j(R) | \frac{\partial^2 \varphi_i(R)}{\partial R^2} \rangle$$

# BOA theoretical analysis

$$\begin{aligned}
 V_{dark} &= \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} \\
 V_{bright} &= \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}
 \end{aligned}
 \rightarrow
 \begin{pmatrix} 0 & 23 & 24 \\ 32 & 0 & 34 \\ 42 & 43 & 0 \end{pmatrix}$$



$$g \begin{pmatrix} 0 & \sin\left(\frac{\theta}{\sqrt{2}}\right) & 2\cos(\theta) \\ \sin\left(\frac{\theta}{\sqrt{2}}\right) & 0 & \sin\left(\frac{\theta}{\sqrt{2}}\right) \\ 2\cos(\theta) & \sin\left(\frac{\theta}{\sqrt{2}}\right) & 0 \end{pmatrix} = g \begin{pmatrix} 0 & \sin(\theta) \\ \sin(\theta) & \cos(\theta) \end{pmatrix} \begin{pmatrix} V_{bright} \\ V_{dark} \end{pmatrix}$$

$$V = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \text{ evolution ?}$$



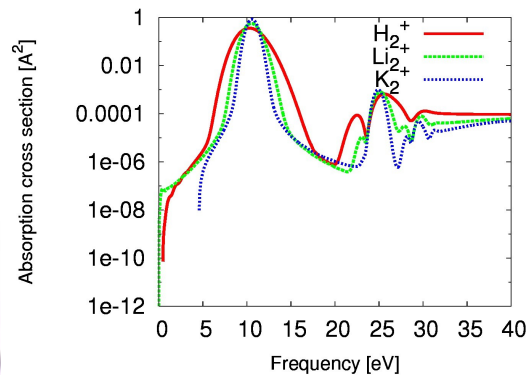
# BOA theoretical analysis

$$i\hbar \frac{\partial v}{\partial t} = g \begin{pmatrix} 0 & \sin(\theta) \\ \sin(\theta) & \cos(\theta) \end{pmatrix} v \longrightarrow v(t)$$

$$c_3(t) = \langle v_{\text{bright}} | v(t) \rangle = \cos^2\left(\frac{\theta}{2}\right) e^{-2ig \sin^2(\frac{\theta}{2})/\hbar} + \sin^2\left(\frac{\theta}{2}\right) e^{2ig \cos^2(\frac{\theta}{2})/\hbar}$$

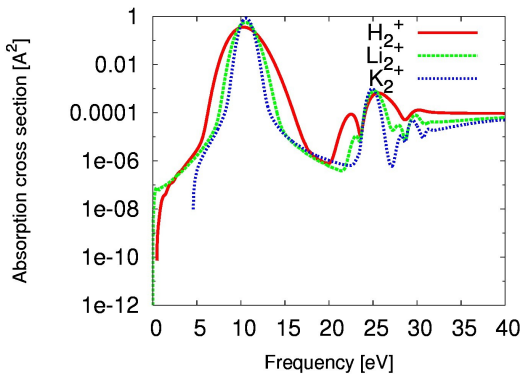
$$d_1(t) = \int \psi_{0(\text{BOA})}^*(R, r, t) r \psi_{1(\text{BOA})}(R, r, t)$$

$$d_3(t) = \int \psi_{0(\text{BOA})}^*(R, r, t) r c_3(t) \psi_{3(\text{BOA})}(R, r, t)$$



# BOA theoretical analysis

$$d(t) = 2k \sum_{i=1,3} d_{i0}^2 \mathfrak{I} \int dR e^{-\frac{E_0 t}{i \hbar}} \chi_0^*(R) c_3(t) \chi_i(R, t)$$



only 3<sup>rd</sup> peak  
peak shift      new 2<sup>nd</sup> peak

$$d(t) = 2k \sum_{i=1,3} d_{i0}^2 \mathfrak{I} \left( e^{-\frac{F_i^2 \sigma_0^2}{2 \hbar^2} t^2 - \left( \omega_i + 2g \sin^2\left(\frac{\theta}{2}\right) - 2g \right) t + \frac{F_i^2}{12 \hbar \mu_I} t^3} \right)$$

$$\sigma_0 = \frac{\hbar^{1/2}}{2^{1/2} \mu_I^{1/4} K^{1/4}}$$

$$\omega_i = \hbar (\varepsilon_i(R_0) - \varepsilon_0(R_0)) - \frac{\hbar K^{1/2}}{4 \mu_I^{1/2}}$$

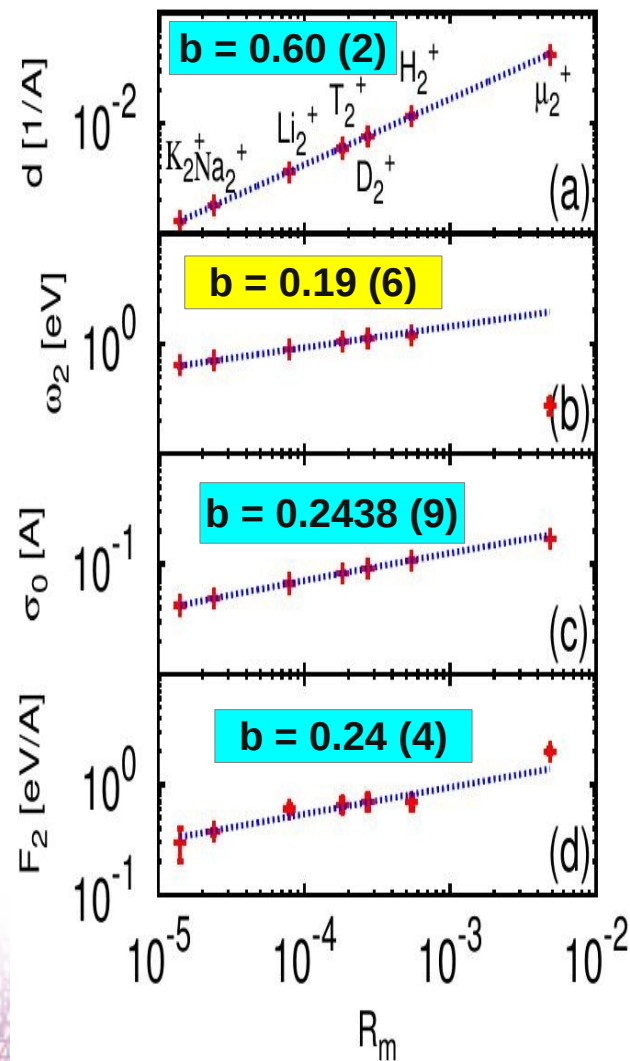
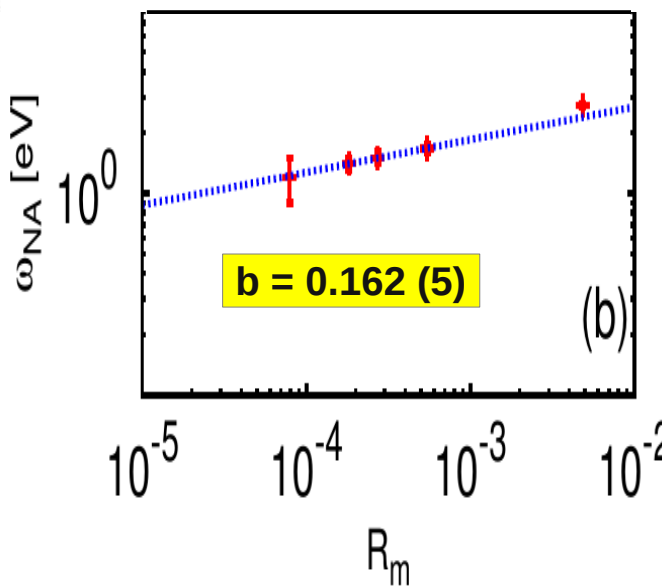
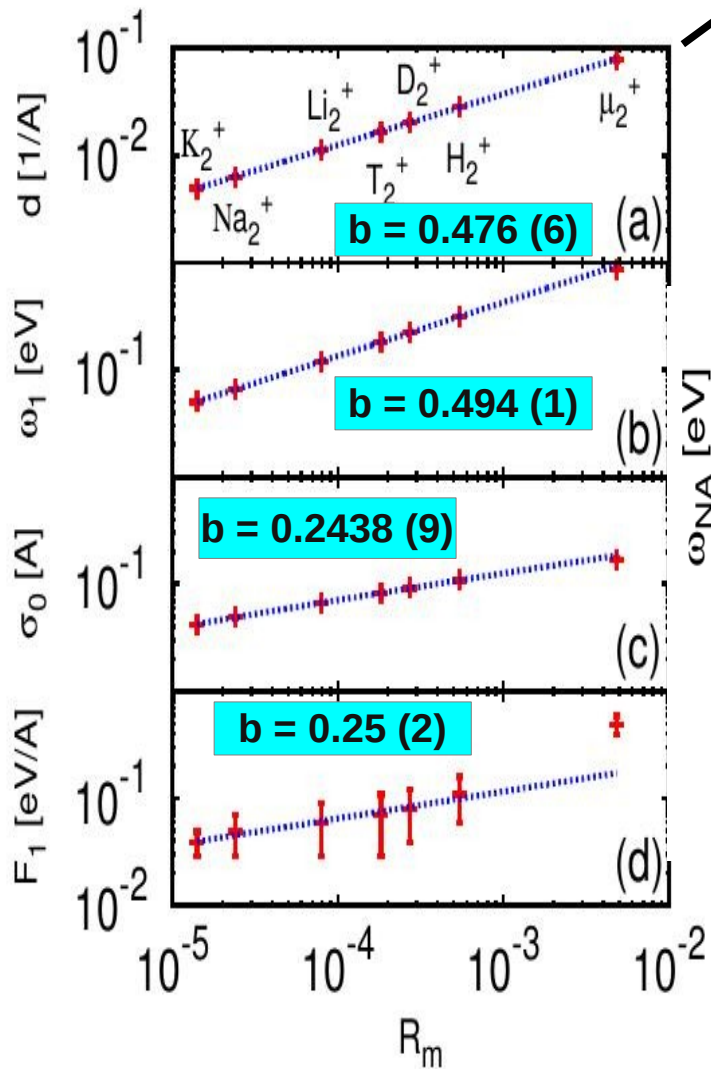
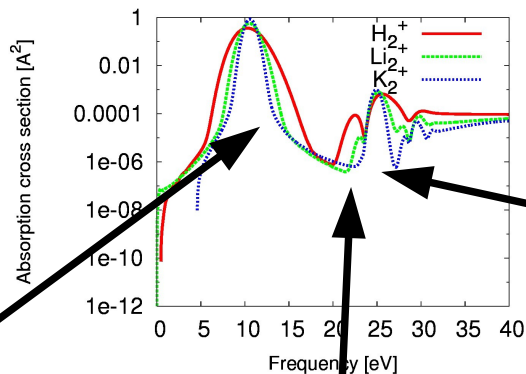
$$F_i = \frac{\hbar^{1/2} K^{3/4}}{2^{1/2} \mu_I^{1/4}}$$

$$d_{i0} = \int \varphi_i^{*(R)}(r) r \varphi_0^{(R)}(r) \rightarrow \varphi(E_0 = \varepsilon_0(R_0) + \frac{1}{2} \hbar \frac{K^{1/2}}{\mu_I^{1/2}})$$



# BOA theoretical analysis

$$y = ax^b$$



# Conclusions

- Non adiabatic effects
  - Gaussian lineshape
  - Peaks shift with ionic mass
  - New non-adiabatic peak
- Theoretically analysed with BOA (fits)
- Electron-ion coupling (BOA not good)
- Other systems ?



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